

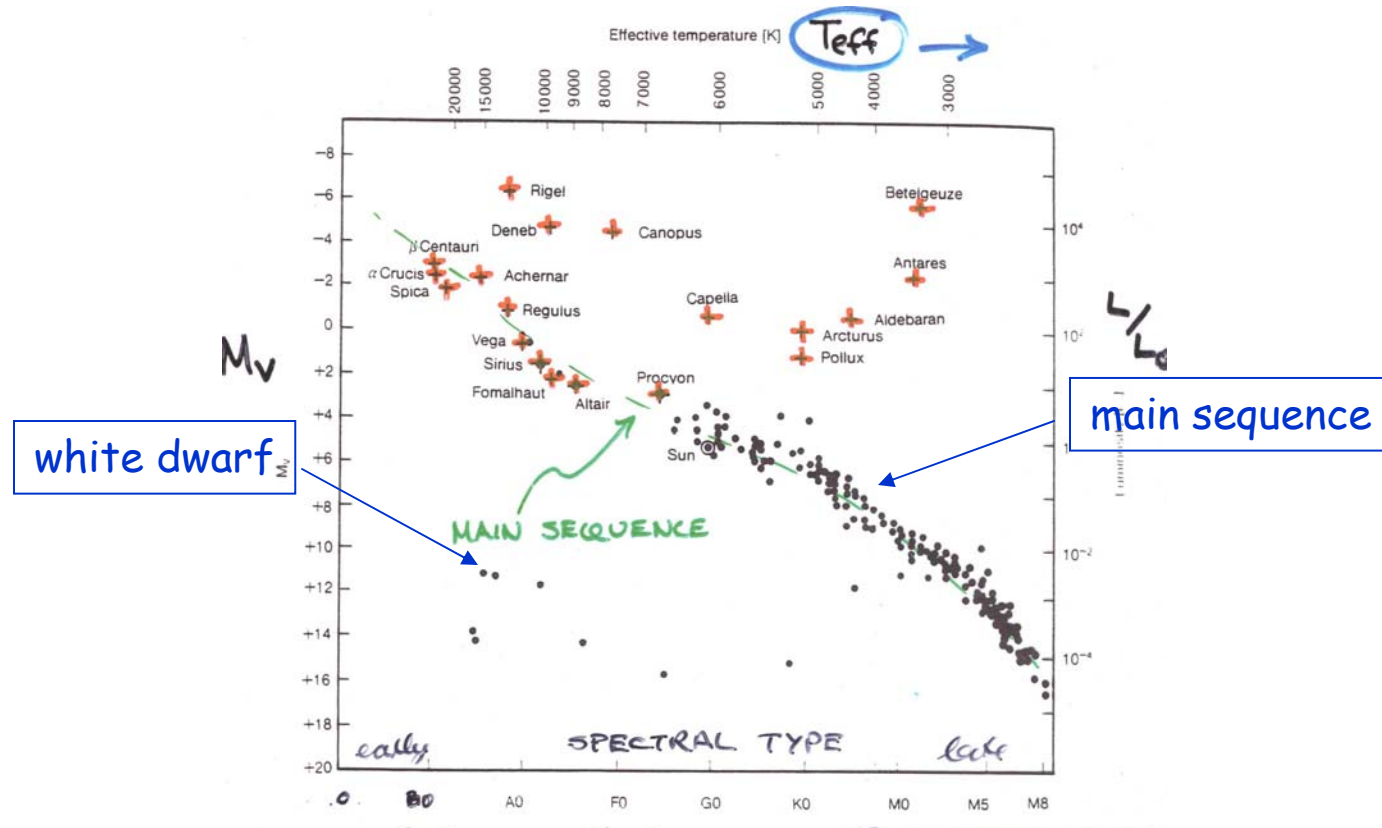
AY 20

Fall 2010

# Hertzsprung-Russell Diagram & Stellar Radiation

Reading: Carroll & Ostlie, Chapter 8.2, Chapter 7

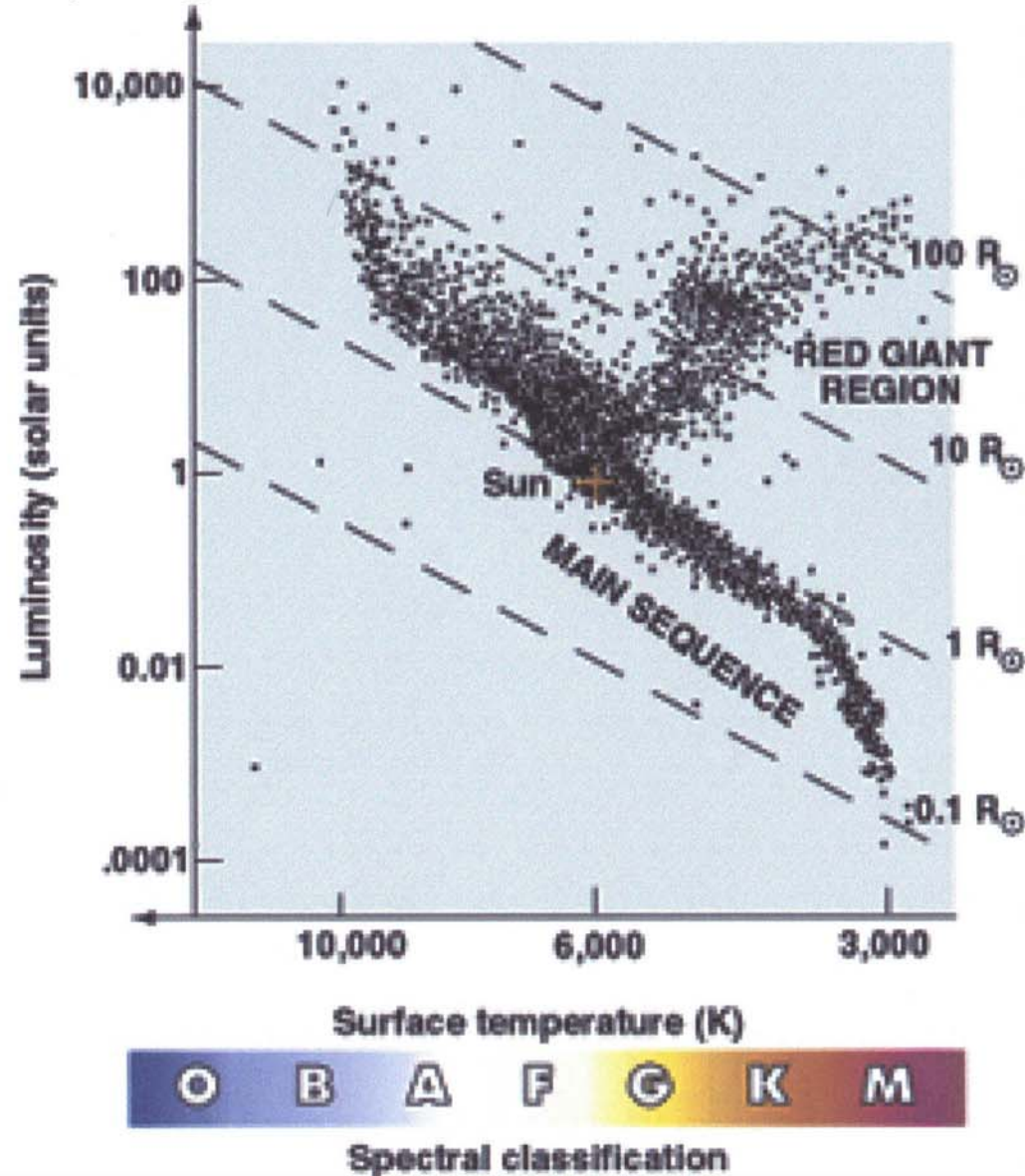
# Hertzprung-Russell Diagram



original form of diagram:  $M_v$  as function of spectral type

- stars within 1 kpc of Sun (solar "neighborhood") + brightest stars (apparent)
- distribution of stars seems to have a pattern; 90% on main sequence
  - HNR → main sequence stars - **dwarfs**; luminous, late spectral type - **giants**
- $M_v$  and spectral type (i.e.  $L$  and  $T_{eff}$ ) are *intrinsic* stellar properties<sub>2</sub>

## H-R (Hertzsprung-Russell) diagram



stars are not spread over  
the entire range of L,T

⇒ clue to the way stars  
work

# Recall Yerkes (MKK) classification system Ia - V

$$L_* = 4\pi R_*^2 F = 4\pi R_*^2 \sigma T_e^4$$

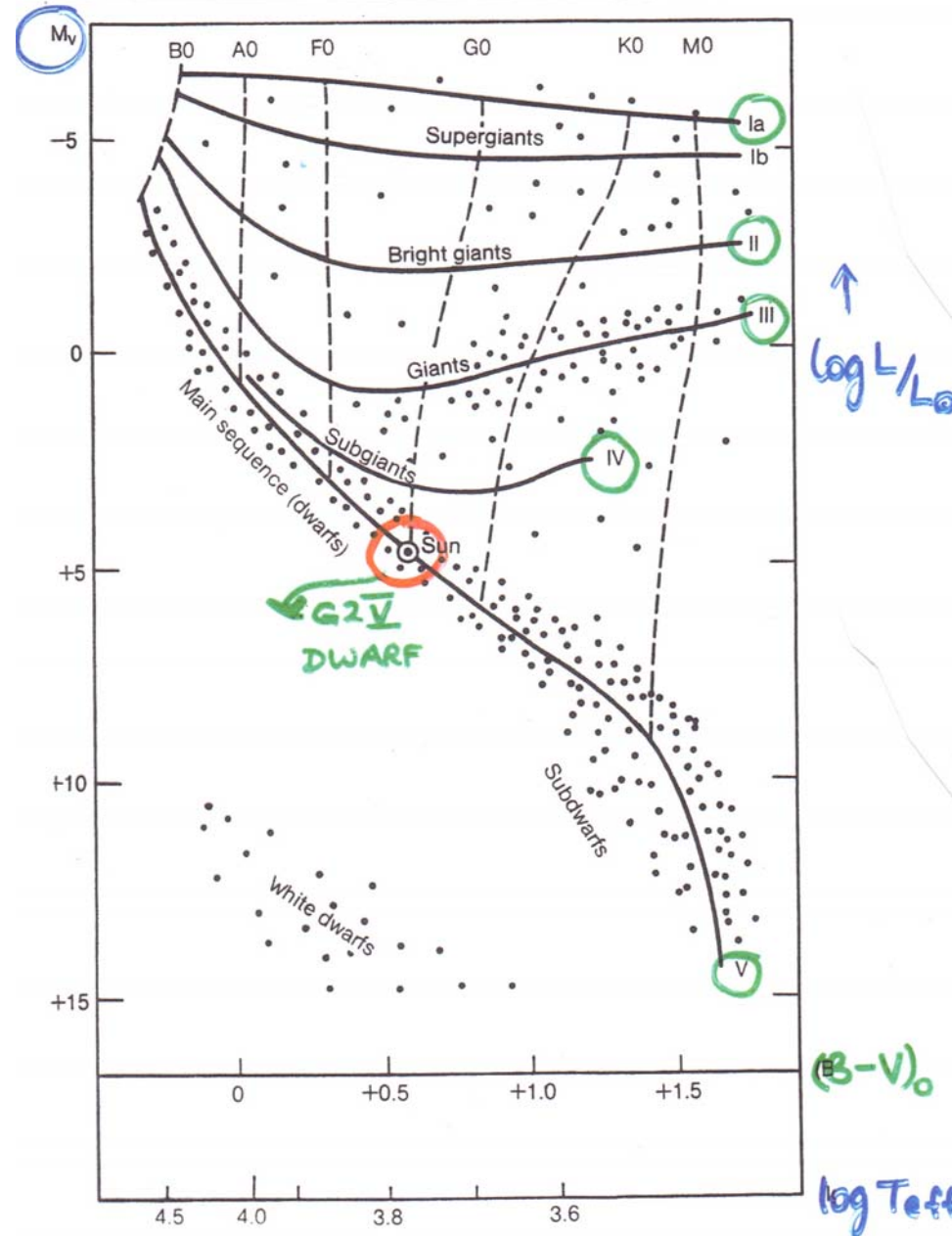
→ absolute magnitude of stars of same spectral type varies with  $R_*$

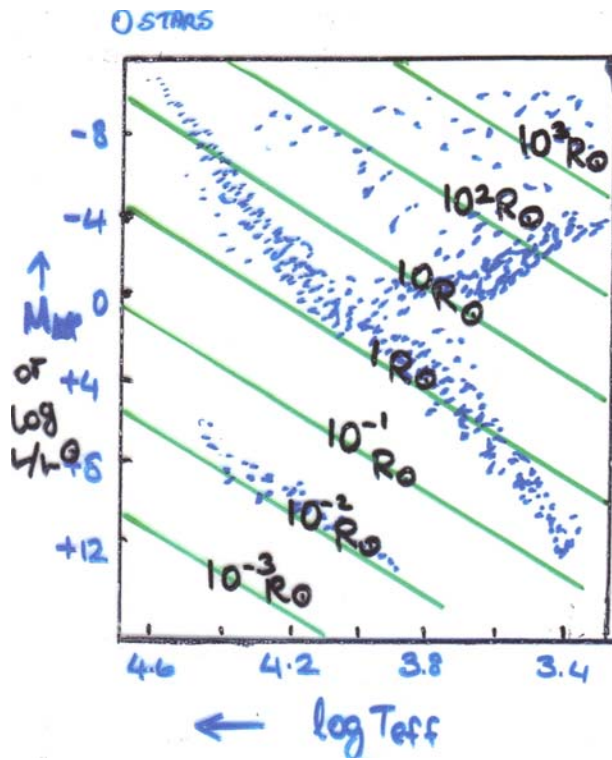
$$R_* = \frac{1}{T_{eff}^2} \sqrt{\frac{L}{4\pi\sigma}}$$

for fixed  $R_*$ ,  $\log L_* \propto \log T_{eff}$

→ lines of constant  $R_*$  in H-R diagram

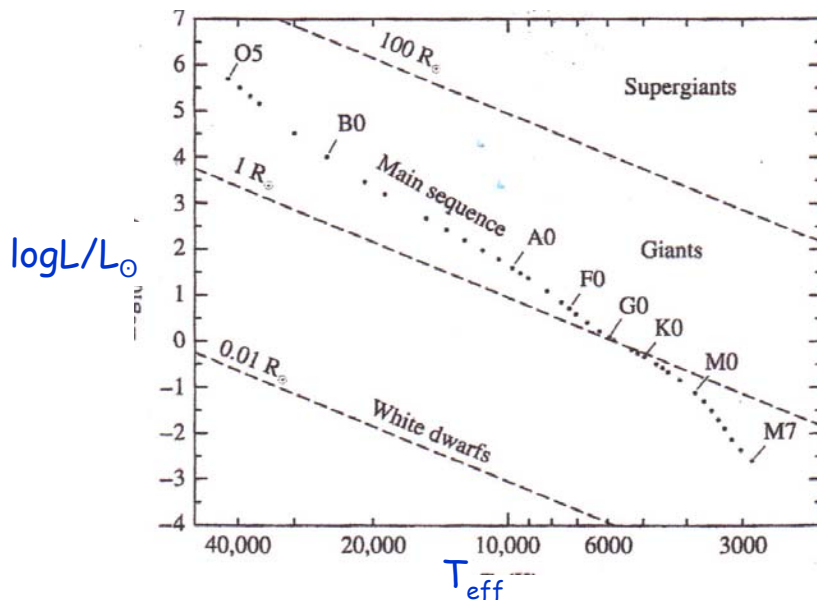
Express  $R_*$ ,  $T_*$ ,  $L_*$ , in terms of  $R_\odot$ ,  $T_\odot$ ,  $L_\odot$ ,





- Supergiants  $R > 100 R_{\odot}$
- Giants  $10 R_{\odot} < R < 100 R_{\odot}$
- main sequence stars (dwarfs)
- $0.1 R_{\odot} < R < 20 R_{\odot}$

• stellar density,  $\rho_* = \frac{M_*}{\frac{4}{3}\pi R_*^3}$   
varies with position in H-R diagram



- C&O 8.2.1: average densities of Sun (G2V), Sirius (A1V), Betelgeuse (~M1a)  
1.4 gms/cm<sup>3</sup> 0.8 gms/cm<sup>3</sup> 10<sup>-8</sup> ρ<sub>⊙</sub>
- Later: position of star on main-sequence depends on its mass



# Stellar Properties from H-R diagram

## 2 The Hertzsprung–Russell Diagram

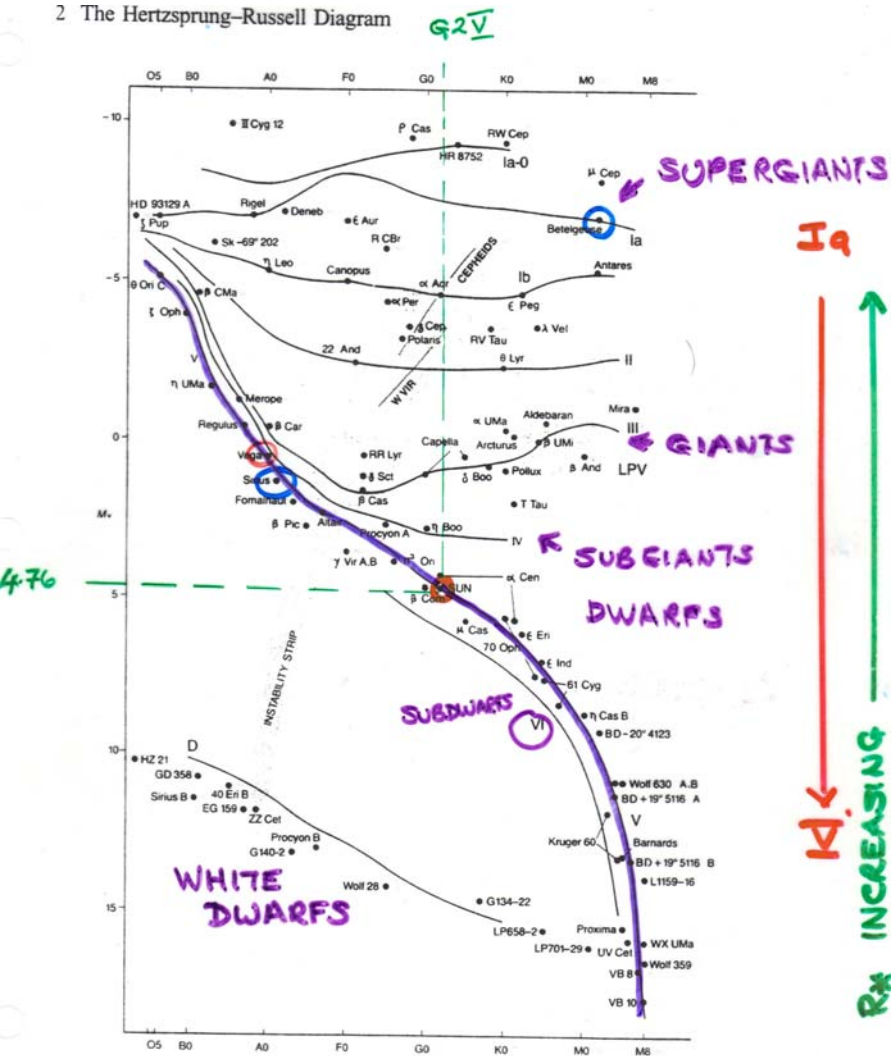


Figure 8.15 Luminosity classes on the H–R diagram. (Figure from Kaler, *Stars and Stellar Spectra*, © Cambridge University Press 1989. Reprinted with the permission of Cambridge University Press.)

spectral type + luminosity class  
(class indicates line width)  $\rightarrow M_v$   
distance from  $m-M = 5 \log d - 5$   
spectroscopic parallax\*

radii from  $L + T_{\text{eff}}$  (see page 5 here),  
hence density if mass known

<u>white dwarfs</u>	<u>supergiants</u>
$R \sim 0.001 R_{\odot}$	$R \sim \text{few} \times 10^3 R_{\odot}$
$\rho \sim 10^9 \text{ kg/m}^3$	$\rho \sim 10^{-4} \text{ kg/m}^3$

(some radii from interferometric measures  
of angular diameters of stars with known  
parallax)

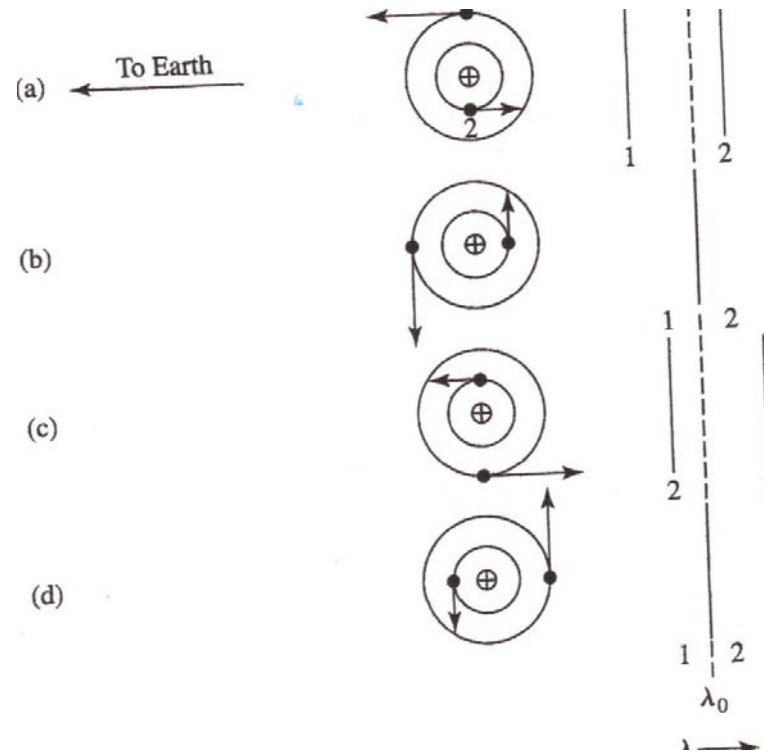
masses from binary star  
observations

# Stellar parameters from binary system measures

- Binary or multiple star systems common
- Classifications
  - Optical doubles - neighbors by coincidence
  - Visual binaries - both members of system can be resolved
    - Can monitor motion if  $P$  not too long → angular sep<sup>n</sup> from ctr of mass
  - Astrometric binaries - only one member visible; its motion reflects presence of a companion
  - Eclipsing binaries - orbital plane oriented so that one star periodically eclipses other → associated variations in light intensity
    - Light curves show two stars → relative  $T_{\text{eff}}$  for each star from depth of minima; radii based on eclipse duration
  - Spectrum binaries - two distinguishable spectra. Orbital period may be so long that no variation of wavelength with time seen

# Spectroscopic binaries - periodic shift in positions of spectral lines is observed

- Lines shift around rest frequency as stars orbit each other
- Easiest to see when velocities in observer's line of sight
- Some component of velocity along line of sight required
- Orbital period should be not too long
- If one component much more luminous than other only one set of shifting lines





# Kepler's Laws generalized for binary systems

Both objects in a binary orbit the center of mass in ellipses with the center of mass at one focus

Kepler's 2<sup>nd</sup> law becomes  $dA/dt = \frac{1}{2} L/\mu = \text{constant}$ ,

$L$  = total angular momentum of system;  $\mu \equiv m_1 m_2 / m_1 + m_2$  = reduced mass  
 $\equiv$  a fixed mass  $\mu$  orbiting center of mass

Integrating for one orbital period:  $t = P$ ,  $A = \frac{1}{2} LP/\mu$

$$A = \pi ab \text{ and } b^2 = a^2(1-e^2), \quad \therefore (LP/2\mu)^2 = \pi^2 a^4(1-e^2)$$

$$\therefore P^2 = 4\pi^2 \mu^2 a^4(1-e^2)/L^2$$

From conservation of angular momentum:

$$L = \mu(GMa(1-e^2))^{1/2}$$

$$\therefore P^2 = 4\pi^2 a^3 / G(m_1 + m_2)$$

$P^2$  inversely proportional to total mass of the system

With  $P$  in years,  $a$  in AU, solar masses, constant = 1

# Stellar masses from visual binaries

For orbital plane perpendicular to line of sight & in center of mass reference frame,\*

$$m_1/m_2 = a_2/a_1$$

$a_2$  and  $a_1$  are semi major axes of 2 ellipses

Angles subtended by these axes a distance  $d$ ,  $\alpha_1$  and  $\alpha_2$

$$\alpha_1 \approx a_1/d \text{ and } \alpha_2 \approx a_2/d$$

$$\therefore m_1/m_2 = \alpha_2/\alpha_1 \rightarrow \text{mass ratio}$$

$$\text{From } P^2 = 4\pi^2 a^3 / G(m_1 + m_2),$$

get  $(m_1 + m_2)$  if semi-major axis of or of reduced mass ( $a=a_1+a_2$ ) known need  $d$ ;

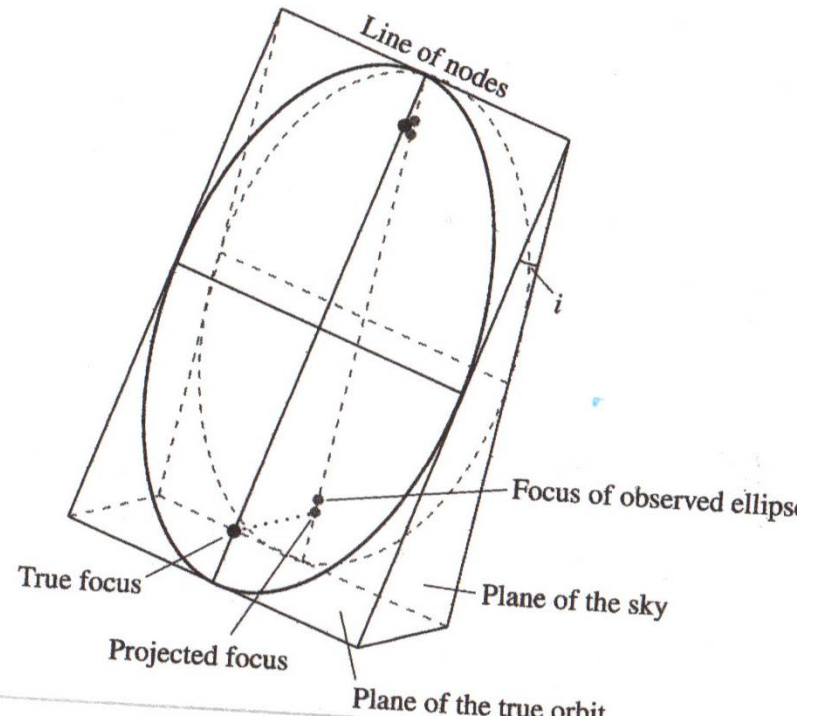
$(m_1 + m_2)$  and  $m_1/m_2 \rightarrow$  masses

May have to take into account proper motion of C of M

Inclination of orbital plane also needed

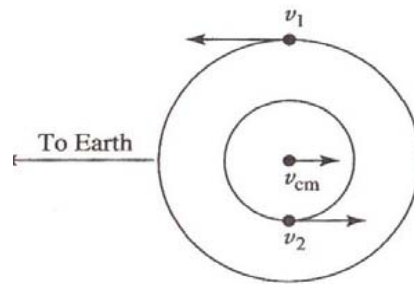
$$m_1 + m_2 = = 4\pi^2 \alpha^3 / G P^2 (d/\cos i)^3$$

: Mass Determination Using Visual Binaries

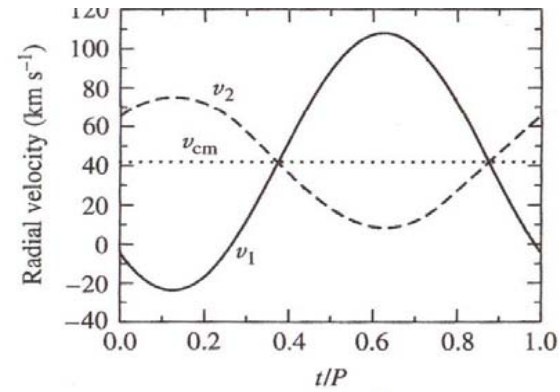


# Masses from double-line spectroscopic binaries

Masses  $1M_{\odot}$  and  $2M_{\odot}$   
 $P = 30\text{d}$ ;  $v_{\text{CofM}} = 42\text{km/s}$



(a)



(b)

Max observed radial velocities  $v_r = v_1 \sin i$ ,  $v_2 \sin i$ ; sinusoidal curves if  $i = 90^\circ$ ;  
 changing  $i$  merely changes amplitude of sinusoids

For  $e \ll 1$ , velocities constant  $v_1 = 2\pi a_1 / P$ ,  $v_2 = 2\pi a_2 / P$

$\therefore m_1 / m_2 = v_2 / v_1$ ; substituting  $\rightarrow m_1 / m_2 = v_{2r} / v_{1r}$

ratio of masses independent of  $\sin i$

From  $a = a_1 + a_2 = P / 2\pi \times (v_1 + v_2)$  and Kepler's 3<sup>rd</sup>:

$$m_1 + m_2 = P / 2\pi G \times (v_1 + v_2)^3$$

$$\therefore m_1 + m_2 = P / 2\pi G \times (v_{1r} + v_{2r})^3 / \sin^3 i$$

But have to be able to measure both radial velocities

# Single line spectroscopic binary

- One star much more luminous than other
- e.g. planets orbiting other stars
- Replace  $v_{2r}$  in terms of stellar masses and  $v_{1r}$

$$m_1 + m_2 = P v_{1r}^3 / 2\pi G \sin^3 i \times (1 + m_1/m_2)^3$$

→ Mass Function Equation

$$\frac{m_2^3 \sin^3 i}{(m_1 + m_2)^2} = \frac{v_1^3 P}{2\pi G}$$

Mass function → only lower limit to mass of  $m_2$   
if  $\sin i$  unknown

# Mass Luminosity Relation

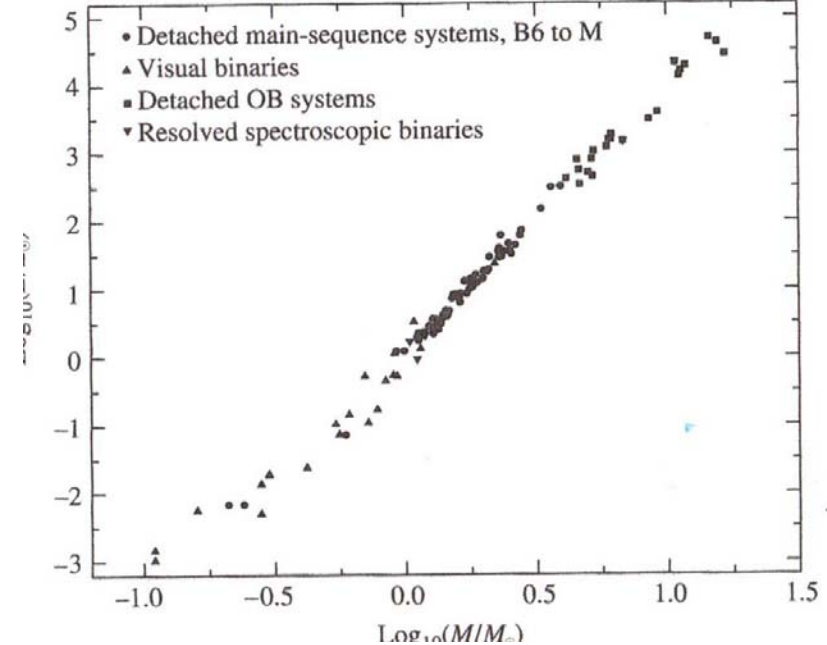
Masses of binary stars  
define empirical  
mass-luminosity relation

For  $M > 3 M_{\odot}$ ,  $L \propto M_{\star}^3$

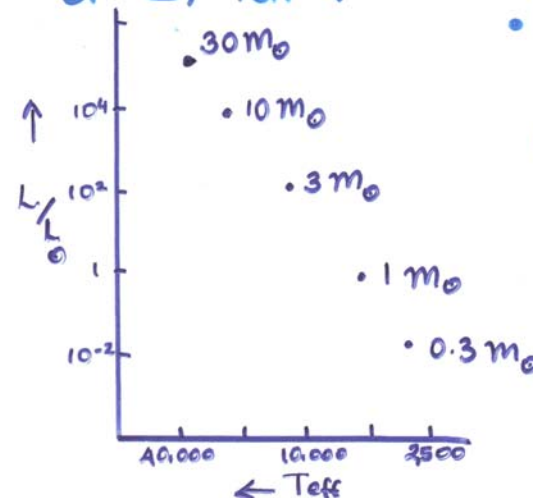
For  $M < 0.5 M_{\odot}$ ,  $L \propto M_{\star}^{2.5}$

VERY approximate

Plotting stars of known mass  
on H-R diagram ( $L$  v.  $T_{\text{eff}}$ )  
also instructive



H-R DIAGRAM IN TERMS  
OF  $L$ ,  $T_{\text{eff}}$  ↓



• PLOTTING  $L$  v  $T_{\text{eff}}$   
FOR KNOWN MAIN  
SEQUENCE BINARIES  
(MASSES DERIVABLE)



EARLY TYPE  
STARS HOTTER,  
& MORE MASSIVE

# Stellar properties from observables

## Direct measurements

- Distances (parallax)
- Luminosities (U,B,V etc)
- Masses (binaries)
- Radii (eclipsing binaries, interferometric measures)

## Spectra also provide info on:

- Rotational velocities
- Chemical abundances
- Magnetic fields
- Mass inflow-outflow

## Using stellar spectra:

- Spectral type  $\equiv T_{\text{eff}}$
- Luminosity classes  $\equiv$  gravity, pressure, density
- Radial velocities,  $z$

## Position on HR Diagram:

- Stellar radii
- Distances (spectroscopic parallax)
- main-sequence masses (approx)
- Ages of stars (later)