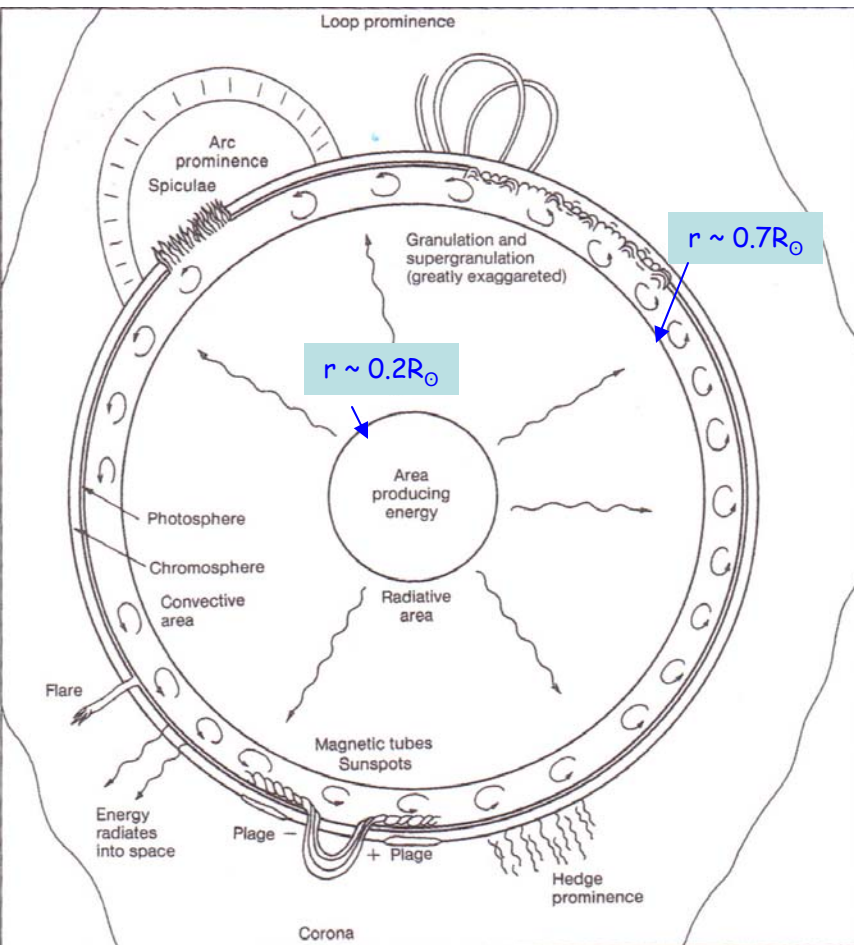


AY 20

Fall 2010

# The Solar Cycle Interstellar Dust

Reading: Carroll & Ostlie, Chapter 11.3, Chapter 12.1



Solar structure  
Reading: C & O §11.2  
"descriptive"

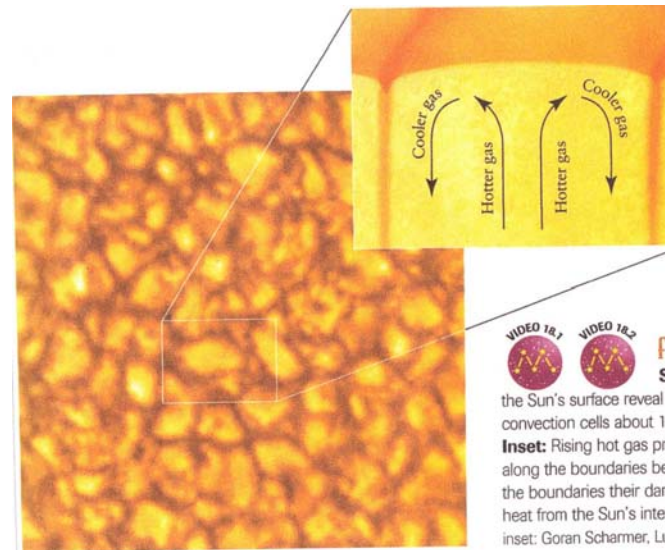
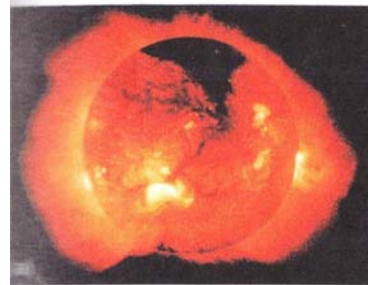
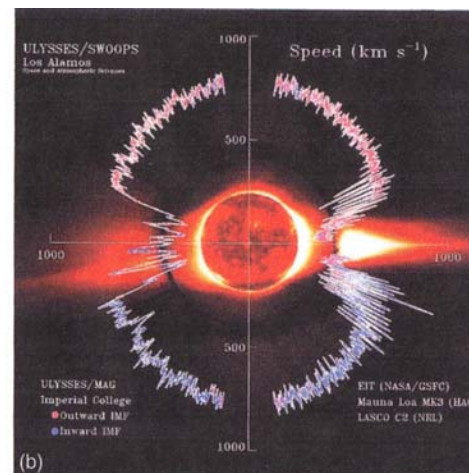


Figure 18-10 R I V U X G

**Solar Granulation** High-resolution photographs of the Sun's surface reveal a blotchy pattern called granulation. Granules are convection cells about 1000 km (600 mi) wide in the Sun's photosphere.  
**Inset:** Rising hot gas produces bright granules. Cooler gas sinks downward along the boundaries between granules; this gas glows less brightly, giving the boundaries their dark appearance. This convective motion transports heat from the Sun's interior outward to the solar atmosphere. (MSFC/NASA inset: Goran Scharmer, Lund Observatory)



**FIGURE 11-24** (a) An X-ray image of the Sun, in false color, taken by the *Yohkoh* satellite. The large, dark area (corresponding to a cooler, less dense region) is a coronal hole. (b) A composite image of the solar wind and corona, in false color, with additional information on the solar magnetic field and speed of the solar wind. Data and images were taken by the *Ulysses* and *SOHO* spacecraft (ESA/NASA missions) and correspond to the 1994 period of sunspot minimum.

# Solar Wind: from expansion of corona (Parker)

Solar corona: high  $T$ , ionized gas  $\equiv$  plasma

High conductivity  $\rightarrow$  isothermal

Hydrostatic equilibrium:  $dP/dr = -GM_{\odot}\rho/r^2$  (here  $M_r \approx M_{\odot}$ )

Completely ionized gas: number density of protons  $n_p \approx \rho/m_p$   
(since  $n_p \approx \rho/m_H$ )

From ideal gas law  $P = 2nkT \therefore d/dr(2nkT) = -GM_{\odot}\rho/r^2$

$$\therefore dn/dr = -GM_{\odot}/2kT \times n_p m_p / r^2$$

Integrating  $\rightarrow n(r) = n_0 e^{-\lambda(1-r_0/r)}$   $\lambda \equiv GM_{\odot}m_p / 2kTr_0$ ;  $n = n_0$  at  $r = r_0$

$$\therefore P(r) = P_0 e^{-\lambda(1-r_0/r)} \text{ and } P_0 = 2n_0 kT$$

limiting values for  $P$  by adopting  $r_0 = 1.4R_{\odot}$ ,  $T = 1.5 \times 10^6 K$ ,  $n_0 = 3 \times 10^{13} \text{ m}^{-3}$

$$\therefore \lambda = 5.5, n(\infty) \approx 1.2 \times 10^{11} \text{ m}^{-3} \text{ and } P(\infty) \approx 5 \times 10^{-6} \text{ N m}^{-2}$$

$$\text{For ISM: } n \sim 3 \times 10^5 \text{ m}^{-3} \text{ and } P \sim 3 \times 10^{-14} \text{ N m}^{-2}$$

$\therefore$  material must be moving outward from Sun = solar wind  
assumption of hydrostatic equilibrium must be invalid

# Solar magnetic field implied by coronal holes (X-ray emission) and sunspots as well as by comets

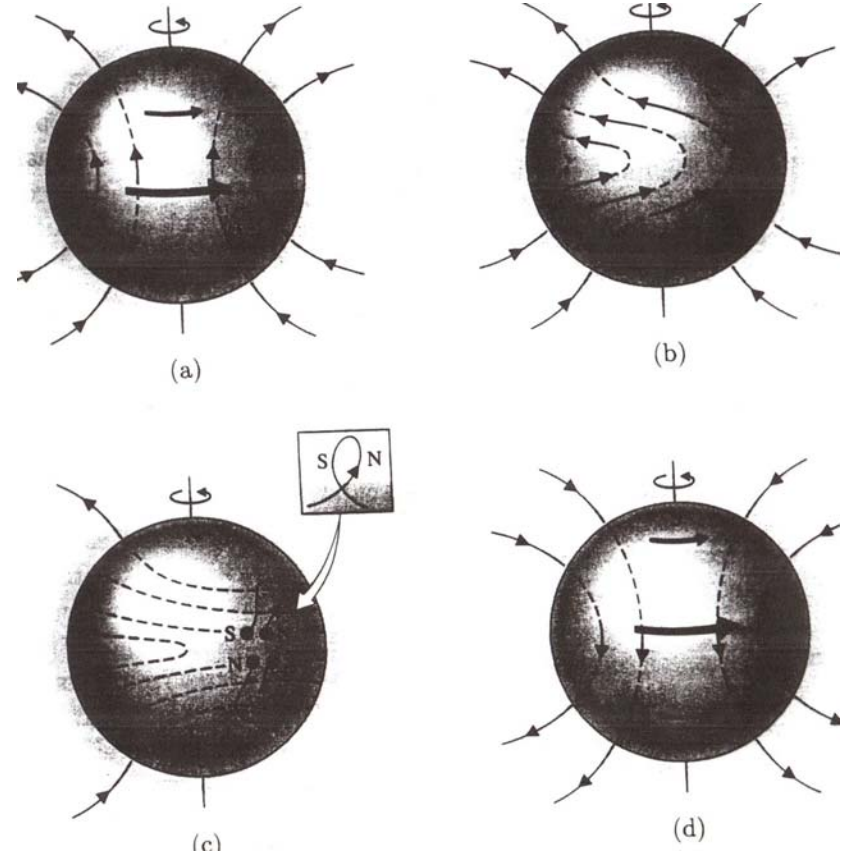
Field due to electric currents in conducting plasma of convection zone (outer 30% of Sun)

Modeled as a magnetic dynamo by Babcock (1961). Generally successful; but detailed MHD treatment required.

Initially field poloidal (a). Field lines "frozen in" gas and dragged by Sun's differential rotation (b); turbulence twists lines into "ropes" - rise to surface as sunspots (c)

Twisting starts at high latitudes, continues to lower. Cancels at equator since polarity of field in lead sunspots  $\equiv$  polarity of original of field (opposite in the two hemispheres)

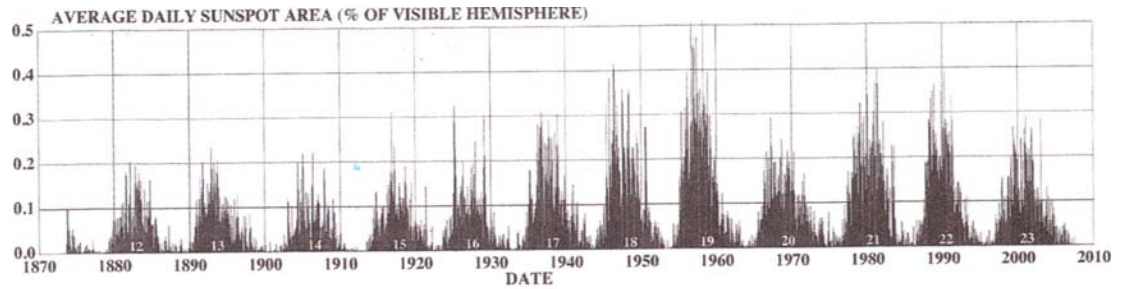
Cycle of migration of groups in 11 yrs. When complete, poloidal field re-established - with *opposite* polarity (d)  $\rightarrow$



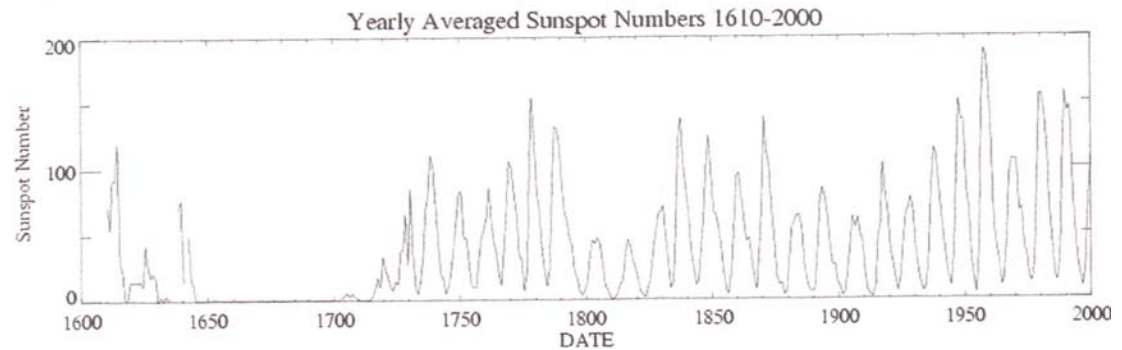
Solar Cycle



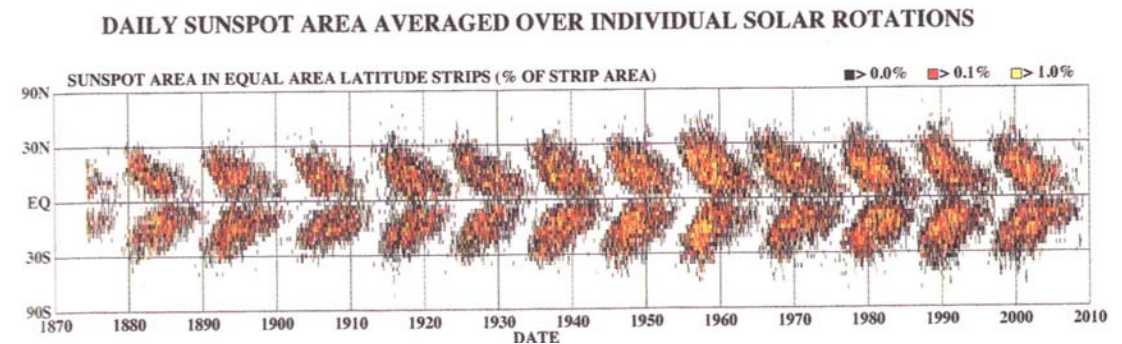
- Variation of number of sunspots with time
- Yearly average number of sunspots. Minimum between 1645 - 1715 = Maunder Minimum
- Little Ice age for Earth
- Butterfly diagram
- Twisting of magnetic field lines Fig 11:37 C&O



(a)

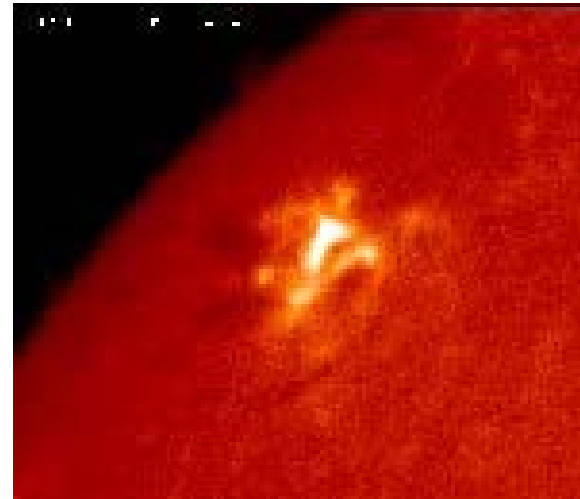


(b)

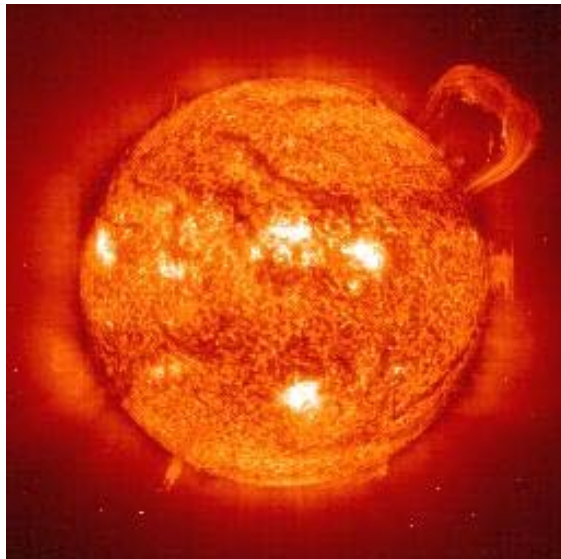




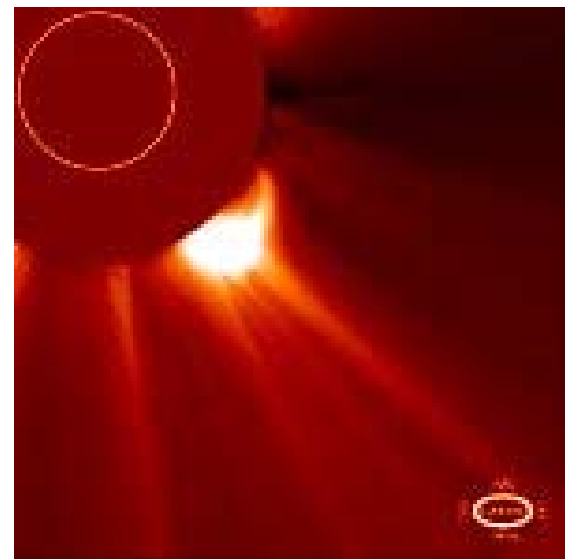
sunspots



flare SOHO



prominences



# Interstellar Medium - the beginning & the end - and everything in between

- ISM -the material - gas and dust - between stars
- Origin of stars
- Processed material from stars returned to ISM via stellar winds, ejection of stellar envelopes, novae, planetary nebulae, supernovae
- Dynamics, structure, and evolution of Milky Way, other galaxies involves ISM
- ISM a laboratory for exotic experiments
- Everything we've learned so far applies: fundamental physics and chemistry, radiative transfer, hydrodynamics, magnetohydrodynamics...
- And it has annoying effects too

# Interstellar Medium & Extinction



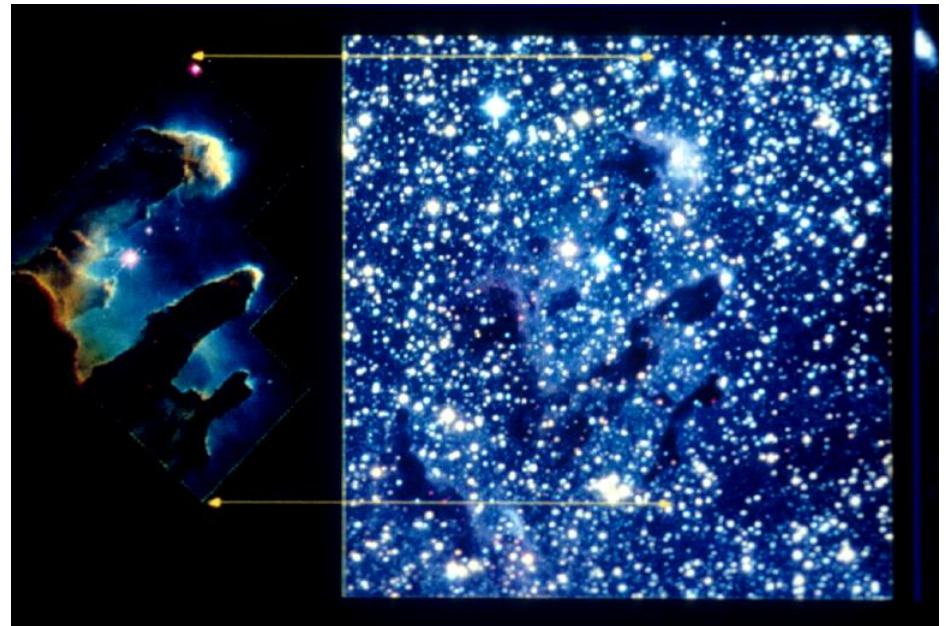
AAT



Hubble



European Southern Observatory ESO



Calar Alto







# Interstellar Extinction

Cluster survey by Trumpler (1930) → star's radiation  $m_\lambda$ , decreased by intervening material (extinction)

$$\text{distance modulus } m_\lambda - M_\lambda = 5 \log d - 5$$

To account for extinction:  $(m_\lambda - A_\lambda) - M_\lambda = 5 \log d - 5$

$$\therefore m_\lambda = M_\lambda + 5 \log d - 5 + A_\lambda \quad (d \text{ in pc, } A_\lambda \text{ extinction in mags})$$

Extinction caused by dust particles = obscuration due to scattering or absorption of radiation ( $\lambda$ -dependent)

scattering  $\equiv$  "reflected"

absorption followed by re-emission at different  $\lambda$ 's

$$\text{Wien Law: } \lambda_{\text{max}} T = 0.29$$

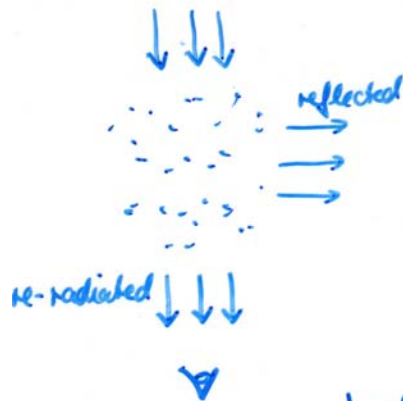
$$\text{ISM: } T \sim 10\text{-}20\text{K} \therefore \lambda_{\text{max}} \sim 300\text{-}150 \mu\text{m}$$

$$\text{near hot star: } T_* \sim 20,000\text{K, } \lambda_{\text{max}} \sim 1450 \text{ \AA}$$

→ UV radiation re-radiated in IR

Average Extinction in plane of Milky Way =  $2^{\text{m}}/\text{kpc}$  11





EXTINCTION DEPENDS

ON  $\lambda$ ,  $\tau$ ,  $n$  WHERE  $n$   
NUMBER  
IS DENSITY OF DUST GRAINS

$$I_{\lambda} = I_{\lambda,0} e^{-\tau_{\lambda}}$$

$$m_1 - m_2 = -2.5 \log(F_1/F_2)$$

Let  $m_{\lambda,0}$  = apparent magnitude real  
 $m_{\lambda}$  = " " " after extinction

$$\Delta m = m_{\lambda} - m_{\lambda,0} = -2.5 \log(F_{\lambda}/F_{\lambda,0}) = -2.5 \log e^{-\tau_{\lambda}} = A_{\lambda}$$

$$\therefore A_{\lambda} = 2.5 \tau_{\lambda} \log e = 1.086 \tau_{\lambda}$$

$\therefore$  OPTICAL DEPTH IN LINE OF SIGHT  $\sim$  CHANGE  
IN MAGNITUDE FROM EXTINCTION

$$\tau_{\lambda} = \int_0^s \kappa_{\lambda} \rho ds \text{ THROUGH 'CLOUD' SIZE } s$$

$I$  FALLS OFF BY  $e^{-1}$  OVER DISTANCE  $l = 1/\kappa_{\lambda} \rho = 1/n_{\lambda}$

$$\sigma_{\lambda} = \text{SCATTERING X-SECTION} \therefore \tau_{\lambda} = \int_0^s n_d(s') \sigma_{\lambda} ds'$$

$n_d(s')$  = NUMBER DENSITY OF SCATTERING GRAINS

$$\therefore \text{FOR CONSTANT } \sigma, \tau_{\lambda} = \sigma_{\lambda} \int_0^s n_d(s') ds' = \sigma_{\lambda} N_d$$

$$N_d = \text{DUST COLUMN DENSITY} \leftarrow \begin{array}{c} \text{area } 1 \text{ m}^2 \\ \text{or } 1 \text{ cm}^2 \end{array} \times$$

$$\therefore A_{\lambda} \sim \tau_{\lambda} \sim N_d (\text{const } \sigma_{\lambda})$$

$\therefore$  EXTINCTION  $\propto$  NUMBER <sup>OF</sup> GRAINS IN LINE OF  
SIGHT



## MIE THEORY FOR DUST PARTICLES:

ASSUME PARTICLES SPHERICAL, RADIUS  $a$

• GEOMETRICAL CROSS-SECTION  $\sigma_g = \pi a^2$

EXTINCTION COEFFICIENT =  $Q_\lambda$  (COMPOSITION DEPENDENT)

$$Q_\lambda = \frac{\sigma_\lambda}{\sigma_g}$$

$\sigma_\lambda$  = SCATTERING CROSS-SECTION

SIMPLISTICALLY, SINCE  $Q_\lambda = \sigma_\lambda / \sigma_g$

Later class: Ay 102:

DEFINE  $x = 2\pi a / \lambda$ ,  $Q_\lambda \Rightarrow$  SERIES EXPANSION IN  $x$

$\Rightarrow$  FOR PARTICLES OF SIZE COMPARABLE TO  $\lambda$

$$x \gtrsim 1 \text{ AND } Q_\lambda \sim a / \lambda$$

$\therefore \sigma_\lambda = Q_\lambda \sigma_g \propto \lambda^{-1}$ , for constant  $a$   $\sigma_\lambda \propto 1/\lambda$  Mie scattering

• IF  $\lambda \gg a$ ,  $Q_\lambda \rightarrow 0$ ; IF  $\lambda \ll a$   $Q_\lambda \rightarrow \text{CONST.}$

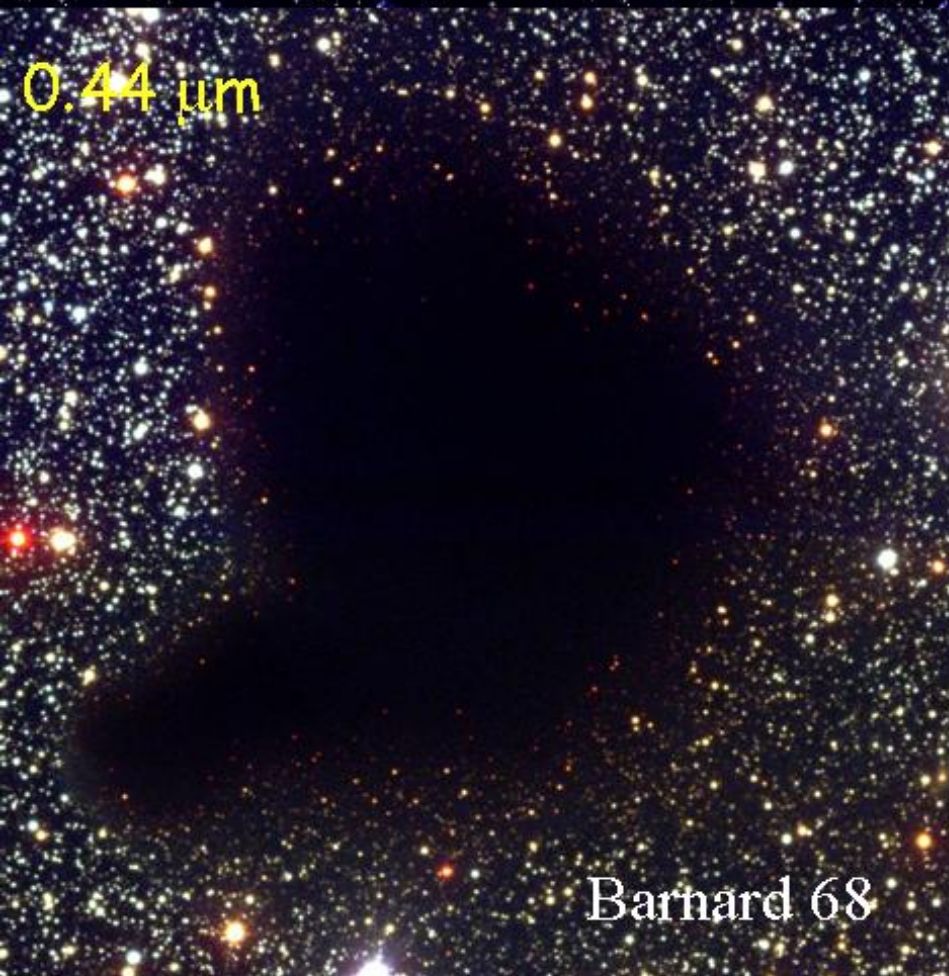
AND  $\sigma_\lambda = \text{const.} a^2 \propto a^2$  Rayleigh Scattering

• EXTINCTION DEPENDENT ON SIZE OF PARTICLES/GRAINS (FOR CONSTANT COMPOSITION)

• EXTINCTION MUCH LESS AT LONG WAVELENGTHS: RADIATION AT IR, MM, CM WAVELENGTHS LESS DIMINISHED

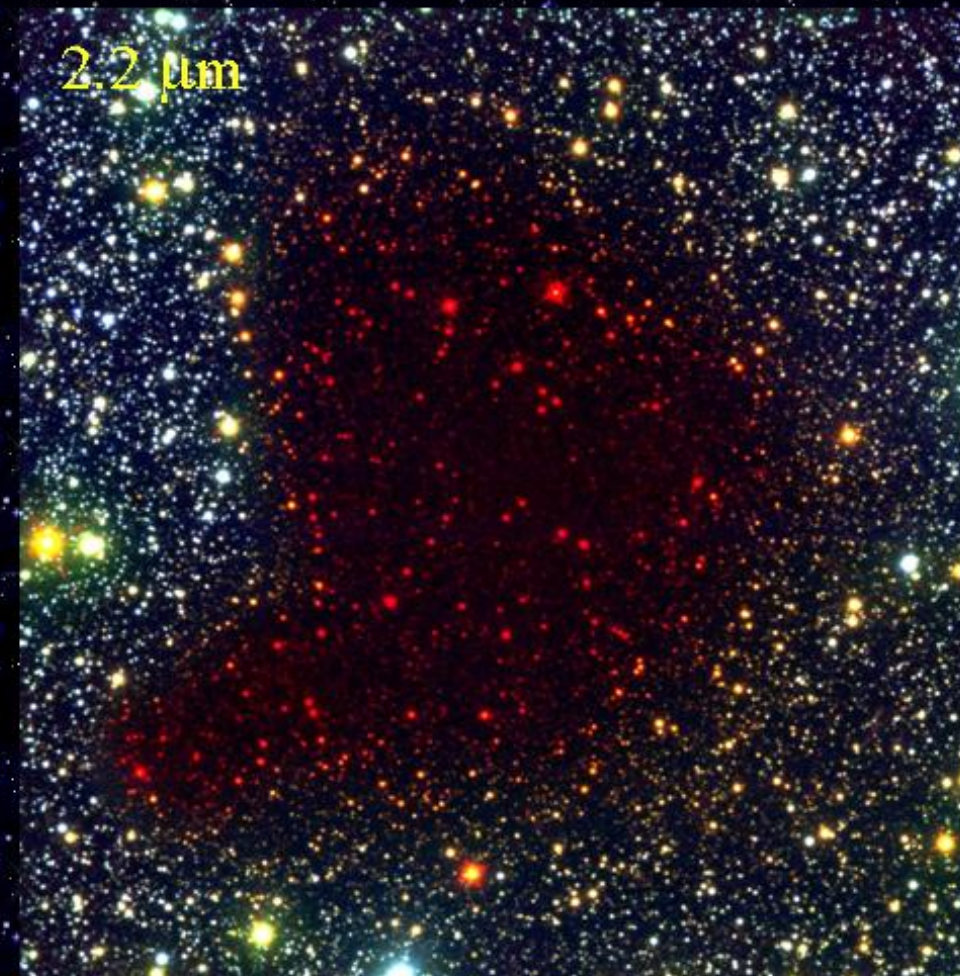


# Less extinction at longer wavelengths



VLT (BVI)

optical image



VLT + NTT (BIR)

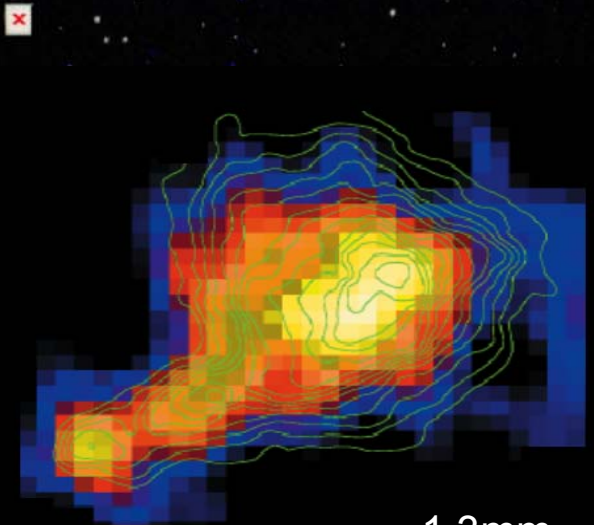
infrared image



# Peering into the darkness

0.44 $\mu\text{m}$

0.55 $\mu\text{m}$

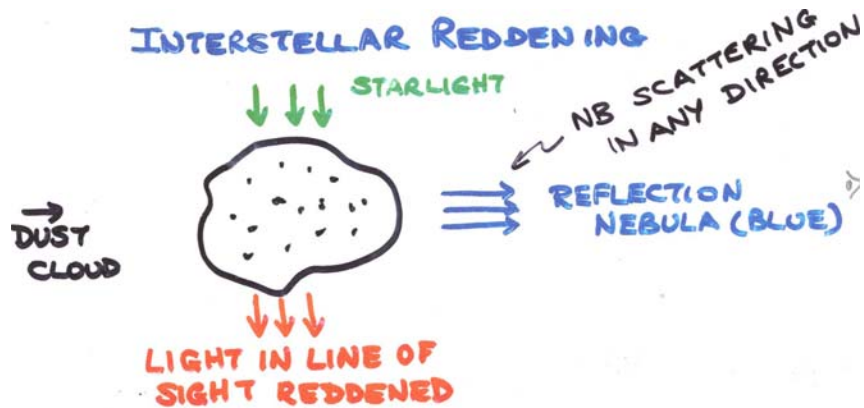


1.3mm

1.25 $\mu\text{m}$

1.65 $\mu\text{m}$

2.16 $\mu\text{m}$



"red"  $\sim 8000\text{\AA}$

"blue"  $\sim 3000\text{\AA}$

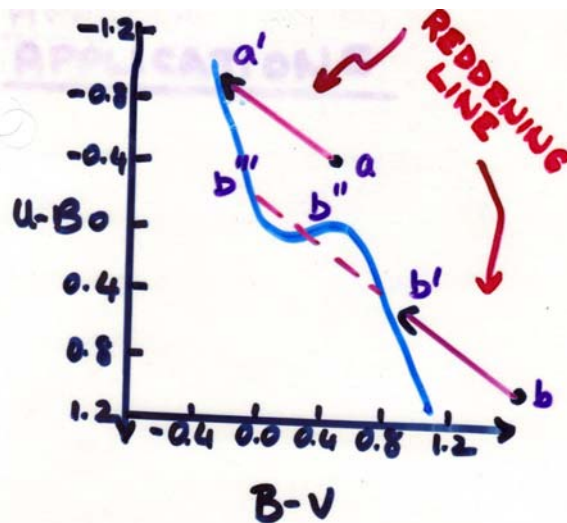
Mie scattering:  $\sigma_\lambda \propto 1/\lambda \therefore$  blue photons scattered more than red  
 $\therefore$  emergent radiation in line of sight redder than when emitted by star  
 AND, if absorbed and re-radiated, emerges at much longer  $\lambda$   
 In presence of dust,  $T_{\text{eff}}$  from observations much lower than at star

$$A_\lambda = 1.086\tau_\lambda \sim \tau_\lambda = N_d, \text{ column density}$$

When  $\lambda \gg a \rightarrow$

Rayleigh scattering  $\sigma_\lambda \propto \lambda^{-4}$  = very strong dependence





- REDDENING DISTORTS DERIVED STELLAR PROPERTIES

e.g. DISTANCE MODULUS,  
 $T_{\text{eff}}$ , SINCE  $\lambda_{\text{max}}$  SHIFTED  
 (WIEN'S LAW)

- CAN 'DE-REDDEN' USING COLOR-COLOR DIAGRAM

REDDENING MEANS MEASURED FLUXES PUT STARS AT POSITION  $a$  RATHER THAN  $a'$

$$\text{COLOR EXCESS} = (\text{B-V})_{\text{MEASURED}} - (\text{B-V})_{\text{INTRINSIC}} = E_{\text{B-V}} = M_{\text{B}} - M_{\text{V}}$$

$$\therefore \text{TRUE } (\text{B-V})_0 = \text{MEASURED } (\text{B-V}) - E_{\text{B-V}}$$

e.g. CASE  $a$ :  $\text{B-V} = 0.4$ ,  $(\text{B-V})_0 = 0$ ,  $E_{\text{B-V}} = -0.4$   
 [FOR CASE  $b$ , SAME EXCESS IF IN SAME CLUSTER ONLY]

$$\text{SINCE } V = M_{\text{V}} + 5 \log d - 5 + A_{\text{V}}, \quad B = M_{\text{B}} + 5 \log d - 5 + A_{\text{B}}$$

$$\text{B-V} = M_{\text{B}} - M_{\text{V}} + A_{\text{B}} - A_{\text{V}} = (\text{B-V})_0 + A_{\text{B}} - A_{\text{V}}$$

$$\therefore E_{\text{B-V}} = (\text{B-V}) - (\text{B-V})_0 = A_{\text{B}} - A_{\text{V}}$$

$$\text{EMPIRICALLY } A_{\text{V}} / E_{\text{B-V}} \approx 3 = R$$

$$\therefore E_{\text{B-V}} = 3 A_{\text{V}} \quad \therefore \text{COLOR EXCESS} \equiv \text{VISUAL EXTINCTION}$$